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*Source / Izvornik:* **Physical Review C - Nuclear Physics, 1998, 58, 1349 - 1351**

**Journal article, Published version**

**Rad u časopisu, Objavljena verzija rada (izdavačev PDF)**

<https://doi.org/10.1103/PhysRevC.58.1349>

*Permanent link / Trajna poveznica:* <https://um.nsk.hr/um:nbn:hr:217:053848>

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*Download date / Datum preuzimanja:* **2024-12-24**



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## Interacting boson-fermion model calculation of the level scheme and electromagnetic properties of $^{99}\text{Zr}$

S. Brant and V. Paar

*Department of Physics, Faculty of Science, University of Zagreb, Zagreb 10000, Croatia*

A. Wolf

*Nuclear Research Center Negev, Beer-Sheva 84190, Israel*

(Received 20 October 1997)

A theoretical study of the  $^{99}\text{Zr}$  nucleus is presented. Energy levels, wave functions, and electromagnetic rates were calculated in the framework of the interacting boson-fermion model and are compared to the available data for low-lying states. In particular, we discuss the sensitivity of the  $g$  factor of the  $3/2_1^+$  state to the quenching of the spin gyromagnetic ratio and to the structure of the respective wave function. [S0556-2813(98)07208-2]

PACS number(s): 27.60.+j, 21.60.Fw, 21.10.Ky

In Ref. [1] the low-lying spherical states and electromagnetic properties of the  $N=59$  nucleus  $^{97}\text{Sr}_{38}$  were described in the framework of the interacting boson-fermion model (IBFM). The purpose of this Brief Report is to present a similar calculation for  $^{99}\text{Zr}_{40}$ , compare it with available experimental data, and in particular with the recently reported result for the  $g$  factor of the  $3/2_1^+$  state [2].

It is well known that the theoretical approach to nuclear structure in the framework of the interacting boson model (IBM) [3], the interacting boson-fermion model (IBFM) [4], and the interacting boson-fermion-fermion model (IBFFM) [5] is capable of treating even-even, odd-even, and odd-odd transitional nuclei, respectively. In this framework the structure of some nuclei in the  $A=100$  region was investigated in recent years [1,6–8]. In particular, two  $N=59$  isotones have been studied so far:  $^{97}\text{Sr}_{38}$  [1] and  $^{98}\text{Y}_{39}$  [7]. In Ref. [1] it was pointed out that the structure of  $^{97}\text{Sr}_{38}$  and  $^{99}\text{Zr}_{40}$  is similar, with the coexistence of spherical and deformed states: the lowest-lying triplet of states  $1/2_1^+, 3/2_1^+, 7/2_1^+$  is of spherical character, while above 0.5 MeV there appear states of deformed nature. These deformed states were associated with the [422]  $3/2$  and [541]  $3/2$  Nilsson orbitals. In  $^{97}\text{Sr}$  the lowest member of the [541]  $3/2$  band was assigned to the  $3/2^-$  level at 644.7 keV and the lowest member of the [422]  $3/2$  band to the  $3/2^+$  level at 585.1 keV. In  $^{99}\text{Zr}$  the possible candidates for these two band heads are the  $3/2^-$  level at 613.96 keV and the  $3/2^+$  level at 724.30 keV [8].

One of the main difficulties encountered in IBFM calculations in cases where sufficient experimental data is not available is the relatively large number of parameters. We approached this problem by considering the similarity of the low-lying states in  $^{99}\text{Zr}$  and  $^{97}\text{Sr}$ . Of all the interaction strengths only the boson-fermion dynamical interaction was adjusted for  $^{99}\text{Zr}$ . The other boson and quasiparticle parameters are taken the same as in the previous IBFM calculation for  $^{97}\text{Sr}$  [1] with one modification: in Ref. [1] it was assumed that the  $\nu d_{5/2}$  configuration was completely occupied and thus it was omitted from the calculation, while in the present calculation for  $^{99}\text{Zr}$  the  $\nu \tilde{d}_{5/2}$  quasiparticle was included, with  $\epsilon(\tilde{d}_{5/2})=1.3$  MeV,  $v^2(\tilde{d}_{5/2})=0.86$ . The values of the boson-fermion interaction strengths  $A_0$  and  $\Lambda_0$  were taken

from the previous calculation for  $^{97}\text{Sr}$  [1]. The dynamical boson-fermion interaction strength  $\Gamma_0$  was changed from the value  $\Gamma_0=0.5$  MeV for  $^{97}\text{Sr}$  to  $\Gamma_0=0.8$  MeV for  $^{99}\text{Zr}$ . This increase of  $\Gamma_0$  causes a lowering of the first two excited states. The value of the quadrupole strength parameter  $\chi$  is taken to be zero in the calculation for  $^{99}\text{Zr}$ , in accordance with the calculation of the energy spectrum.

In Fig. 1 the calculated IBFM states of spherical type in  $^{99}\text{Zr}$  are presented in comparison to the available experimental data. As in the case of the  $N=59$  isotone  $^{97}\text{Sr}$ , the wave function of the  $1/2^+$  ground state of  $^{99}\text{Zr}$  is dominated by the  $\tilde{s}_{1/2}$  quasiparticle, while the  $3/2^+$  first excited state is of a more complex character. Components of the latter wave function larger than 1% are given in Table I, expressed in the boson-fermion coupled basis:

$$|3/2_1^+\rangle = \sum_{jn_d\nu I} \xi_{jn_d\nu I}^{3/2} |j, n_d\nu I; \frac{3}{2}\rangle. \quad (1)$$

In the basis state  $|j, n_d\nu I; \frac{3}{2}\rangle$  the quasiparticle  $j$  and the  $n_d$   $d$ -boson state  $|n_d\nu I\rangle$  of angular momentum  $I$  are coupled to the total angular momentum  $\frac{3}{2}$ . Here,  $|n_d\nu I\rangle$  denotes the IBM basis state  $|n_d\nu I, n_s=N-n_d; I\rangle$ , where  $n_d$   $d$  bosons and  $n_s$   $s$  bosons are coupled to the total angular momentum  $I$ . The quantity  $\nu$  denotes an additional quantum number, if needed, which distinguishes the  $d$ -boson states having the same values of quantum numbers  $n_d I$ . As seen from Table I, the wave function of the  $3/2_1^+$  state does not have a single dominant component. The two largest components in the  $|3/2_1^+\rangle$  wave function are comparable: the quasiparticle state  $|\tilde{d}_{3/2}, 00; 3/2\rangle$  (27.8%) and the one- $d$ -boson multiplet state based on the  $\tilde{g}_{7/2}$  quasiparticle  $|\tilde{g}_{7/2}, 12; 3/2\rangle$  (27.4%). The fact that two components account for more than 50% of the wave function is significant: it means that, within the limitations of the present model, the state has a predominantly spherical character—as opposed to deformed states, whose wave functions are expected to contain many different configurations, with no particular preference for any of them. A possible way to check this statement is by comparing experimental values of observables of the  $3/2_1^+$  state to calculations



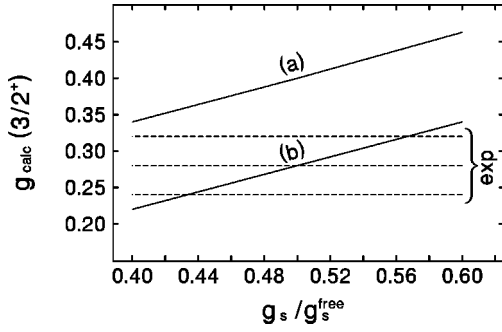


FIG. 2. Calculated  $g$  factor of the  $3/2_1^+$  state versus the quenching factor  $g_s/g_s^{\text{free}}$ . The solid curves present results for: (a)  $g_T = -2.02$ ; (b)  $g_T = 0$ . The dashed lines mark the experimental  $g$  factor with its error bars.

overall interval. Furthermore, we note that this value of the quenching factor  $g_s/g_s^{\text{free}}$  is similar to the one used in previous calculations for several nuclei in the  $A = 100$  region, like  $^{103}\text{Ru}$  and  $^{105}\text{Pd}$  [6].

The value chosen here for the tensor interaction strength  $g_T$  is somewhat smaller than the value used in the previous calculation for  $^{97}\text{Sr}$  [1]. We also note that the value  $g_T \approx +0.5g_s^{\text{free}}$  is rather close to an overall estimate of  $g_T \approx +0.4g_s^{\text{free}}$  from Ref. [11].

In Table II, the calculated electromagnetic properties corresponding to the above parametrization are compared to the available data under the previously mentioned assumption for the spin and parity of the 575.4 keV level. The agreement between the calculations and experiment for the few observables for which we have data is within a factor of 2–3, which can be considered as reasonable and in fact typical for IBFM calculations of electromagnetic observables in medium and heavy nuclei. Now we discuss in some detail the structure of the  $3/2_1^+$  wave function and the value of its magnetic moment. The experimental value of the  $g$  factor is  $g(3/2_1^+) = +0.28(4)\mu_N$  [2]. For the parametrization in Eq. (2) we obtain in IBFM the value  $g(3/2_1^+) = +0.34$ . On the other hand, without the contribution from the tensor term in the  $M1$  operator the calculated result is  $g(3/2_1^+) = +0.22$ . Moreover, the calculated value of  $g(3/2_1^+)$  is sizably dependent on the value of the quenching factor  $g_s/g_s^{\text{free}}$ . This is shown in Fig. 2 by the solid line labeled (a), by keeping the values of the

TABLE III. Contributions of various components of the wave function to the  $g$  factor of the  $3/2_1^+$  state of  $^{99}\text{Zr}$ .

$j$	$n_d$	$I$	$j'$	$n'_d$	$I'$	Contribution to $g(3/2_1^+)$
5/2	2	4	5/2	2	4	0.011
7/2	1	2	7/2	1	2	0.027
7/2	2	2	7/2	2	2	0.010
7/2	2	4	7/2	2	4	0.018
7/2	3	4	7/2	3	4	0.006
1/2	1	2	1/2	1	2	0.041
1/2	1	2	3/2	1	2	0.011
3/2	0	0	3/2	0	0	0.130
3/2	1	2	1/2	1	2	0.011
3/2	1	2	3/2	1	2	0.043
3/2	2	0	3/2	2	0	0.020
3/2	2	2	3/2	2	2	0.010
3/2	3	0	3/2	3	0	0.007

other parameters ( $g_R, g_l, g_T$ ) fixed as given by the parametrization in Eq. (2). The solid line labeled (b) presents the value of  $g(3/2_1^+)$  versus  $g_s/g_s^{\text{free}}$ , without inclusion of the tensor term. From Fig. 2 it is clear that at present, one cannot unambiguously determine the importance of the tensor term in the  $M1$  operator, although it is clear that for  $g_T = 0$  a better agreement with the experimental value is obtained for the entire range of  $g_s/g_s^{\text{free}}$  values.

In Table III we present the contributions to the calculated  $g(3/2_1^+)$  from various components of the wave function. We see that about 70% of the experimental value of the  $g$  factor is due to three matrix elements, thus supporting the above contention that the  $3/2_1^+$  state has a predominantly spherical character.

In conclusion, the IBFM provides a reasonable description of low-lying states of the transitional  $^{99}\text{Zr}$  nucleus. A detailed comparison of the calculated value of the magnetic moment of the first excited state with a recently reported experimental value supports the expectation that this state has spherical character, in accordance with the conclusions of previous works regarding shape coexistence in this nucleus.

We wish to thank Dr. Gerard Lhersonneau for useful comments and discussions.

[1] G. Lhersonneau, B. Pfeiffer, K.-L. Kratz, H. Ohm, K. Sistemich, S. Brant, and V. Paar, *Z. Phys. A* **337**, 149 (1990).  
 [2] A. Wolf, R. L. Gill, Z. Berant, and D. S. Brenner, *Phys. Rev. C* **51**, 2381 (1995).  
 [3] A. Arima and F. Iachello, *Phys. Rev. Lett.* **35**, 1069 (1975).  
 [4] F. Iachello and O. Scholten, *Phys. Rev. Lett.* **43**, 679 (1979).  
 [5] V. Paar, in *Capture Gamma-Ray Spectroscopy and Related Topics*, edited by S. Raman, AIP Conf. Proc. No. 125 (AIP, New York, 1984), p. 70; S. Brant, V. Paar, and D. Vretenar, *Z. Phys. A* **319**, 351 (1984).  
 [6] S. Brant *et al.*, in *Proceedings of the International Workshop on Nuclear Structure of the Zirconium Region*, edited by J.

Eberth, R.A. Meyer, and K. Sistemich (Springer-Verlag, Berlin, 1988), p. 199.  
 [7] S. Brant, V. Paar, G. Lhersonneau, O. W. B. Schult, H. Seyfarth, and K. Sistemich, *Z. Phys. A* **334**, 517 (1989).  
 [8] S. Brant, K. Sistemich, V. Paar, and G. Lhersonneau, *Z. Phys. A* **330**, 365 (1988).  
 [9] A. Arima and H. Horie, *Prog. Theor. Phys.* **12**, 623 (1954).  
 [10] R. J. Blin-Stoyle and M. A. Perkes, *Proc. R. Soc. London, Ser. A* **67**, 885 (1954).  
 [11] A. Bohr and B. R. Mottelson, *Nuclear Structure* (Benjamin, New York, 1969), Vol. I.