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## Measurement of the $^{240}\text{Pu}(n,f)$ cross-section at the CERN n\_TOF facility: First results from experimental area II (EAR-2)

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**Abstract.** The accurate knowledge of the neutron-induced fission cross-sections of actinides and other isotopes involved in the nuclear fuel cycle is essential for the design of advanced nuclear systems, such as Generation-IV nuclear reactors. Such experimental data can also provide the necessary feedback for the adjustment of nuclear model parameters used in the evaluation process, resulting in the further development of nuclear fission models. In the present work, the  $^{240}\text{Pu}(\text{n},\text{f})$  cross-section was measured at CERN's n\_TOF facility relative to the well-known  $^{235}\text{U}(\text{n},\text{f})$  cross section, over a wide range of neutron energies, from meV to almost MeV, using the time-of-flight technique and a set-up based on Micromegas detectors. This measurement was the first experiment to be performed at n\_TOF's new experimental area (EAR-2), which offers a significantly higher neutron flux compared to the already existing experimental area (EAR-1). Preliminary results as well as the experimental procedure, including a description of the facility and the data handling and analysis, are presented.

## 1. Introduction

The optimum design and development of advanced nuclear systems requires the accurate knowledge of neutron induced fission cross sections of plutonium isotopes and other minor actinides [1,2]. More specifically, the  $^{240}\text{Pu}$  isotope is included in the Nuclear Energy Agency's (NEA) High Priority Request List (HPRL) [3] as well as in the NEMEA-4 meeting proceedings [4]. Requested target accuracies are in the order of 1.6%–11.8% for the energy region 0.454 keV–6.07 MeV compared to the presently reported 3.9%–21.6%. The previous experiment carried out at the n\_TOF's Experimental Area 1 (EAR-1), was unsuccessful due to detector deterioration caused by the high alpha activity of the  $^{240}\text{Pu}$  samples in addition to the long period needed to complete the measurement [5], therefore the neutron induced fission cross section had to be remeasured in the newly constructed experimental area EAR-2, which offers a higher background suppression and a significantly higher instantaneous neutron flux, thus reducing the time needed to complete the measurement. Similar recent measurements have also been performed at GNEIS [6], LANSCE [7] and IRMM [8].

## 2. Experimental setup

### 2.1. The n\_TOF facility

The measurement was performed at CERN's n\_TOF facility [9–11], in the newly commissioned experimental area EAR-2 [12]. The neutrons are produced by spallation, occurring when a 20 GeV/c pulsed proton beam, delivered by CERN's Proton Synchrotron [13], impinges on a 40 cm in length and 60 cm in diameter lead block, resulting in a neutron energy spectrum that spans from the meV to the GeV region. The proton pulse has a duration of 6 ns RMS, a maximum repetition rate of 1 pulse per 1.2 s and nominally carries  $7 \times 10^{12}$  protons. The 1.3 ton lead target is surrounded by a 1 cm-thick layer of water as well as

an additional 4 cm-thick layer of borated water ( $\text{H}_2\text{O} + 1.28\%\text{H}_3\text{BO}_3$ , fraction in mass) which act as coolant and neutron moderator respectively. Following the target in the beam direction, a beam pipe of 182.3 m leads to the first experimental area EAR-1. The new experimental area EAR-2 is located 18.16 m above the spallation target, perpendicular to the beam direction, offering a 25 – 30 times higher flux [14] and approximately 10 times shorter time of flights compared to the existing area. The high instantaneous flux, combined with the wide white neutron spectrum (from thermal energies to the MeV region), leads to the relative suppression of the contribution originating from the high intrinsic  $\alpha$ –background of the samples, thus allowing for accurate measurements below the fission threshold, where the fission cross section is low.

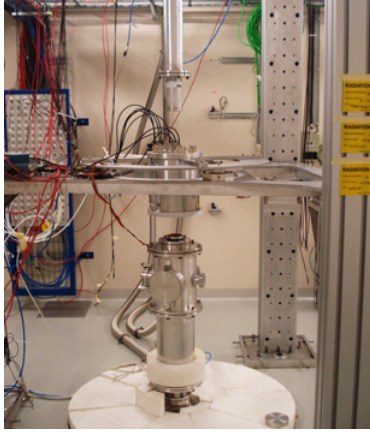
### 2.2. Samples

Three high purity (99.89% atomic abundance) plutonium dioxide ( $\text{PuO}_2$ ) samples, prepared at IRMM (Belgium) [15] for the previous measurement in EAR-1, were used having a total  $^{240}\text{Pu}$  mass of 2.288 mg ( $\sim 100 \mu\text{g}/\text{cm}^2$  per sample) and a total  $\alpha$ -activity of 19.219 MBq. The oxide material was electro-deposited on a 250  $\mu\text{m}$ -thick aluminum backing with a diameter of 5 cm, while the deposit itself had a diameter of 3 cm. In addition to the plutonium samples, two targets were used as reference: a  $^{235}\text{U}$  and a  $^{238}\text{U}$ , with a mass of  $\sim 0.6$  mg and  $\sim 0.8$  mg respectively. The impurities found in the samples, mainly  $^{239}\text{Pu}$ , have a non-negligible contribution in the obtained yields below 1 keV, which was thoroughly taken into account in the analysis.

### 2.3. Detectors

The measurement was carried out using a set-up based on the extremely compact, low mass Micromegas (Micro-Mesh Gaseous Structure) gaseous detector [16–18]. Its volume is divided into two parts by a thin, 35  $\mu\text{m}$  pitched micro-mesh: a drift region (5 mm) and a narrow amplification region (50  $\mu\text{m}$ ), in which a high

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**Figure 1.** The fission chamber, which housed the detectors and samples, as seen inside the experimental area EAR-2, in the middle of the picture. The neutron beam was delivered vertically from bottom to top.

electric field causes an avalanche multiplication. Typical operating electric fields are  $\sim 1$  kV/cm and  $\sim 100$  kV/cm respectively depending on the gas, which in this case was a mixture of  $Ar:CF_4:isoC_4H_{10}$  (88:10:2) at atmospheric pressure and room temperature. The detector set-up consisted of 6 similar detectors in total (3 for the plutonium samples, 2 for the uranium samples and 1 for monitoring possible proton recoils from the surrounding materials), all of which were housed in a cylindrical aluminum chamber as shown in Fig. 1.

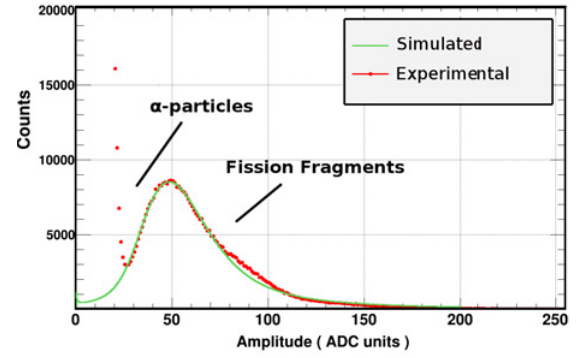
#### 2.4. Data acquisition

The analogue detector signals were recorded using a high performance digital acquisition system [19], operating at a 500 MHz sampling rate with a resolution of 8 bits. To minimize the amount of data recorded during acquisition, a zero-suppression algorithm was applied.

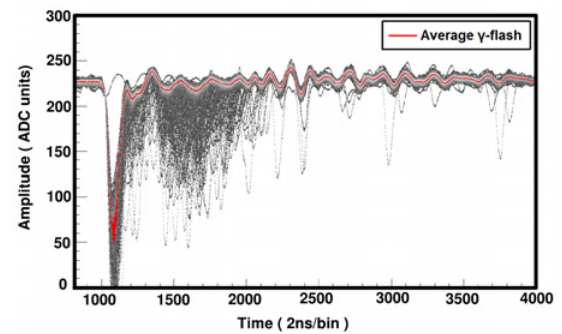
### 3. Data analysis

An offline pulse shape analysis routine [20] parses the raw data, determining, among others, the amplitude and the position in time of every recorded signal. During the experiment, beam-off data was obtained in order to measure the contribution of the  $\alpha$ -activity as well as spontaneous fission. Monte Carlo simulations were also performed, using the general purpose code FLUKA [21,22] to study the behavior of the detectors. More specifically, the pulse height spectra of  $\alpha$ -particles and fission fragments were reconstructed for all detectors, making possible both the estimation of the detection efficiency and the determination of the correction for the signal amplitude threshold, which was found to be in the order of 10%. A typical experimental pulse height spectrum along with the simulated one, is shown in Fig. 2.

The proton interaction with the lead target, apart from the neutron production, results in the generation of prompt  $\gamma$ -rays and relativistic particles, that reach the experimental hall at (nearly) the speed of light. This burst of particles is commonly referred to in literature as “ $\gamma$ -flash”. The  $\gamma$ -flash results in the formation of a high amplitude pulse, comparable to the fission signals, which



**Figure 2.** Experimental (red points) and simulated (green line) typical pulse height spectra for  $^{240}Pu$ . The direct simulated output is calibrated and convoluted with the detector's response function.



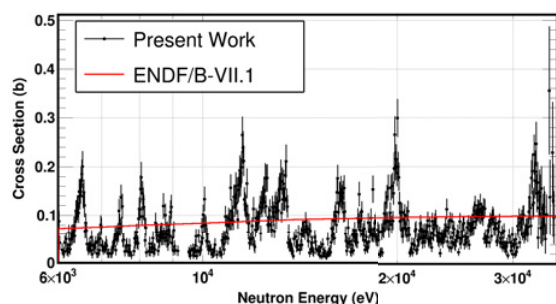
**Figure 3.** Typical average  $\gamma$ -flash (red) and raw  $\gamma$ -flash pulses (gray scale) for  $^{240}Pu$ . During the extraction of the average, the fission signals are rejected. The average  $\gamma$ -flash shape is typically calculated from one thousand raw signals.

lasts a few hundreds of ns and is followed by a baseline oscillation in times of flight that correspond to energies in the MeV region. To mitigate this effect, an average  $\gamma$ -flash shape is derived for every detector, which is subtracted individually from each  $\gamma$ -flash pulse, during the offline raw data processing. A typical average  $\gamma$ -flash is shown in Fig. 3.

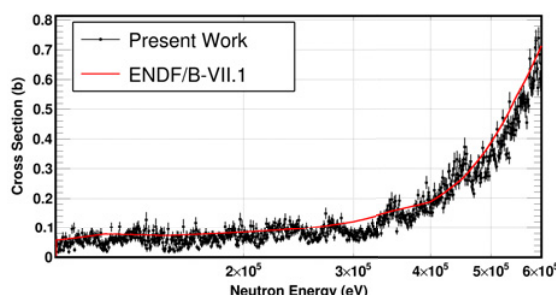
The  $^{240}Pu(n, f)$  cross section was calculated using the reference reactions  $^{235}U(n, f)$  and  $^{238}U(n, f)$  for flux normalisation purposes. The detector dead time, as well as the background contribution due to sample impurities were taken into account during the analysis. The final experimental cross section was derived from the weighted average between the three individual cross sections, calculated for each plutonium target. Preliminary results are shown in Fig. 4 for the energy range 6 keV–35 keV, while resonances can be resolved up to a few tens of keV. The present preliminary data comparison with the latest ENDF/B-VII.1 library [23,24], reveals that the aforementioned resonances are not included in the evaluations, as Fig. 4 demonstrates. Resonance analysis is planned to be performed, using R-matrix calculations, to characterize the resolved resonances.

In the energy range above 35 keV, data has been obtained up to at least several MeV. Specifically in the sub-threshold region between 100 keV and 600 keV, the shape of the extracted cross section is in general agreement with the ENDF/B-VII.1 evaluations, as shown in Fig. 5.





**Figure 4.** Comparison between present data (black points) and ENDF/B-VII.1 evaluation (red line) in the energy range 6 keV–35 keV. Clusters of resonances, that are attributed to transitions between Class-I and Class-II states are visible, while resonances can be resolved up to a few tens of keV. Data is shown using a binning of 2000 bins per energy decade.



**Figure 5.** Comparison between present data (black points) and ENDF/B-VII.1 evaluation (red line) in the energy range 100 keV–600 keV. Data is shown using a binning of 2000 bins per energy decade.

Resonance-like structures are also observed, which are attributed to high transition states. Above 1 MeV, the non optimized sample mass caused a high counting rate which affects the normalisation, therefore analysis is still ongoing.

## 4. Conclusion

The  $^{240}\text{Pu}$  neutron induced fission cross section, is the first measurement to be performed in the newly commissioned experimental area EAR-2, at CERN's n\_TOF facility. Data was obtained over a wide energy range that spans from thermal energies up to at least a few MeV. Most remarkably, data in the sub-threshold region, shows clear resonance structures which are not included in the evaluated libraries. R-matrix calculations will be performed to characterize these structures. In the MeV region, preliminary data shows a general agreement with

the evaluations, yet further analysis is ongoing. The higher instantaneous flux along with the stronger background suppression, compared to the existing experimental area EAR-1, are deemed crucial to the success of this measurement, following which additional measurements are planned that include the study of short-lived and rare isotopes such as  $^{230}\text{Th}$ ,  $^{232}\text{U}$ ,  $^{238,241}\text{Pu}$  and  $^{244}\text{Cm}$ .

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